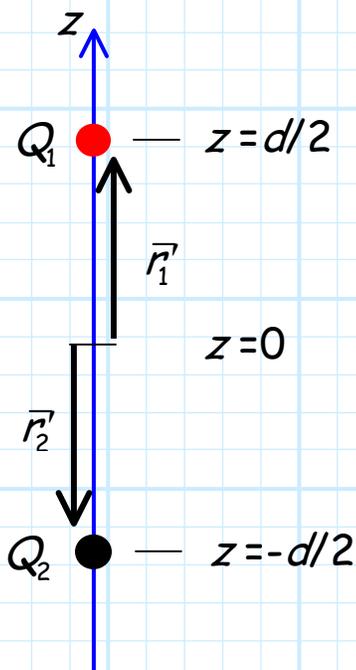


# Example: The Electric Dipole

Consider two point charges ( $Q_1$  and  $Q_2$ ), each with **equal magnitude** but **opposite sign**, i.e.:

$$Q_1 = Q \quad \text{and} \quad Q_2 = -Q \quad \text{so} \quad Q_1 = -Q_2$$

Say these two charges are located on the  $z$ -axis, and separated by a **distance  $d$** .



The **location** of charge  $Q_1 = -Q$  is therefore specified by **position vector**  $\vec{r}_1' = \frac{d}{2} \hat{a}_z$

The **location** of charge  $Q_2 = -Q$  is therefore specified by **position vector**  $\vec{r}_2' = \frac{-d}{2} \hat{a}_z$

**The Electric Dipole**

We call this charge configuration an **electric dipole**. Note the **total charge** in a dipole is **zero** (i.e.,  $Q_1 + Q_2 = Q - Q = 0$ ). But, since the charges are located at different positions, the electric field that is created is **not zero**!

**Q:** *Just what is the electric field created by an electric dipole?*

**A:** One approach is to use **Coulomb's Law**, and add the resulting electric **vector** fields from each charge together.

However, let's try a different approach. Let's find the **electric potential field** resulting from an electric dipole. We can then take the gradient to find the electric field!

Note that this should be relatively **straightforward**! We already know the electric potential resulting from a **single** point charge—the electric potential resulting from two point charges is simply the **summation** of each:

$$V(\bar{r}) = V_1(\bar{r}) + V_2(\bar{r})$$

where the electric potential  $V_1(\bar{r})$ , created by charge  $Q_1$ , is:

$$V_1(\bar{r}) = \frac{Q_1}{4\pi\epsilon_0 |\bar{r} - \bar{r}_1|} = \frac{Q}{4\pi\epsilon_0 \left| \bar{r} - \frac{d}{2} \hat{a}_z \right|}$$

and electric potential  $V_2(\bar{r})$ , created by charge  $Q_2$ , is:

$$V_2(\bar{r}) = \frac{Q_2}{4\pi\epsilon_0 |\bar{r} - \bar{r}_2|} = \frac{-Q}{4\pi\epsilon_0 \left| \bar{r} + \frac{d}{2} \hat{a}_z \right|}$$

Therefore the **total** electric potential field is:

$$\begin{aligned} V(\bar{r}) &= \frac{Q}{4\pi\epsilon_0 \left| \bar{r} - \frac{d}{2} \hat{a}_z \right|} - \frac{Q}{4\pi\epsilon_0 \left| \bar{r} + \frac{d}{2} \hat{a}_z \right|} \\ &= \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{\left| \bar{r} - \frac{d}{2} \hat{a}_z \right|} - \frac{1}{\left| \bar{r} + \frac{d}{2} \hat{a}_z \right|} \right) \end{aligned}$$

If the point denoted by  $\bar{r}$  is a significant distance away from the electric dipole (i.e.,  $|\bar{r}| \gg d$ ), we can use the following **approximations**:

$$\frac{1}{\left| \bar{r} - \frac{d}{2} \hat{a}_z \right|} \approx \frac{1}{|\bar{r}|} + \frac{d \cos \theta}{2|\bar{r}|} = \frac{1}{r} + \frac{d \cos \theta}{2r^2}$$

$$\frac{1}{\left| \bar{r} + \frac{d}{2} \hat{a}_z \right|} \approx \frac{1}{|\bar{r}|} - \frac{d \cos \theta}{2|\bar{r}|} = \frac{1}{r} - \frac{d \cos \theta}{2r^2}$$

where  $r$  and  $\theta$  are the **spherical coordinate** variables of the point denoted by  $\bar{r}$ .

Therefore, we find:

$$\begin{aligned}
 V(\bar{r}) &= \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{\left| \bar{r} - \frac{d}{2} \hat{a}_z \right|} - \frac{1}{\left| \bar{r} + \frac{d}{2} \hat{a}_z \right|} \right) \\
 &= \frac{Q}{4\pi\epsilon_0} \left( \left( \frac{1}{r} + \frac{d \cos\theta}{2r^2} \right) - \left( \frac{1}{r} - \frac{d \cos\theta}{2r^2} \right) \right) \\
 &= \frac{Q}{4\pi\epsilon_0} \frac{d \cos\theta}{r^2}
 \end{aligned}$$

Note the result. The **electric potential field** produced by an **electric dipole**, when centered at the **origin** and aligned with the **z-axis** is:

$$V(\bar{r}) = \frac{Qd \cos\theta}{4\pi\epsilon_0 r^2}$$

**Q:** *But the original question was, what is the **electric field** produced by an electric dipole?*

**A:** Easily determined! Just take the **gradient** of the electric potential function, and multiply by -1.

$$\begin{aligned}
 \mathbf{E}(\bar{r}) &= -\nabla V(\bar{r}) \\
 &= -\nabla \left( \frac{Qd \cos\theta}{4\pi\epsilon_0 r^2} \right) \\
 &= \frac{-Qd}{4\pi\epsilon_0} \left[ \cos\theta \frac{d}{dr} \left( \frac{1}{r^2} \right) \hat{a}_r + \frac{1}{r^3} \frac{d(\cos\theta)}{d\theta} \hat{a}_\theta \right] \\
 &= \frac{-Qd}{4\pi\epsilon_0} \left[ \left( \frac{-2 \cos\theta}{r^3} \right) \hat{a}_r - \frac{\sin\theta}{r^3} \hat{a}_\theta \right]
 \end{aligned}$$

The static **electric field** produced by an **electric dipole**, when centered at the **origin** and aligned with the **z-axis** is:

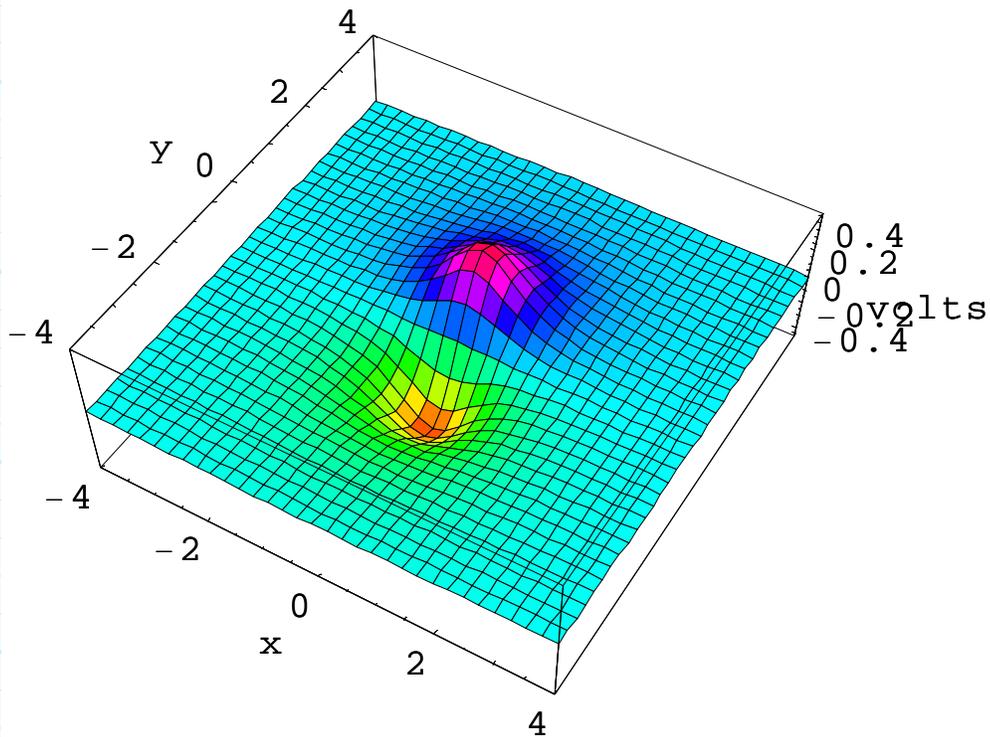
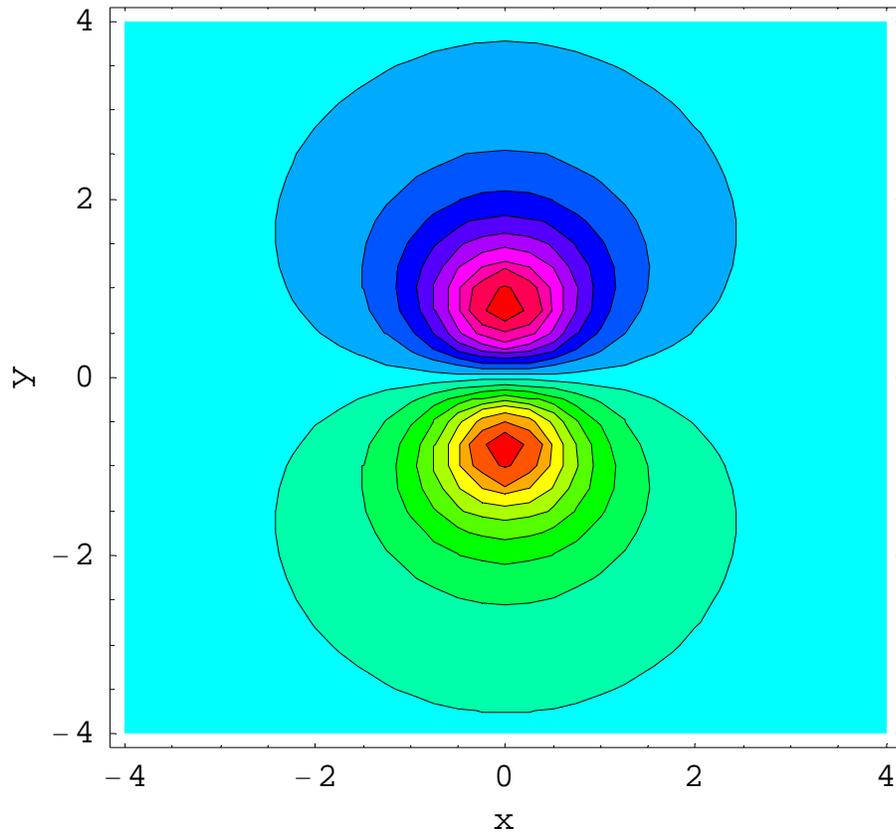
$$\mathbf{E}(\bar{r}) = \frac{Qd}{4\pi\epsilon_0} \frac{1}{r^3} \left[ 2\cos\theta \hat{a}_r + \sin\theta \hat{a}_\theta \right]$$

Yikes! **Contrast** this with the electric field of a **single** point charge. The electric dipole produces an electric field that:

- 1) Is proportional to  $r^{-3}$  (as opposed to  $r^{-2}$ ).
- 2) Has vector components in **both** the  $\hat{a}_r$  and  $\hat{a}_\theta$  directions (as opposed to just  $\hat{a}_r$ ).

In other words, the electric field does **not** point away from the electric dipole!

The **electric potential** produced by an electric dipole looks like:



And the **electric field** produced by the electric dipole is:

